# TOWARDS A NEW MEASUREMENT OF THE ELECTRON EDM

# COLE MEISENHELDER CUA PIZZA TALK



# **SECTION I: WHY ARE EDMS INTERESTING?**

### MOTIVATION

- The Standard Model cannot answer all questions
- Dark matter and dark energy remain unexplained
- Matter-Antimatter
  Asymmetry Problem



## **POSSIBLE SOLUTION TO BARYON ASYMMETRY**

- Sakharov Conditions (1967) allow for Baryon Asymmetry
- CP Violation
  - Allowed in quark mixing
  - Standard Model does not account for the observed asymmetry
  - Need physics Beyond the Standard Model (BSM)
  - CPT theorem makes CP and T violations equivalent



# EDMS AND CP VIOLATION

- Permanent EDMs inherently violate P and T
- Standard Model predicts electron EDM  $< 10^{-38} e \cdot cm$
- Current limit:  $|d_e| < 1.1 \times 10^{-29} e \text{ cm}$
- Many BSM theories predict nonzero eEDMs near our sensitivity



# PLACING LIMITS ON NEW PHYSICS



SUSY particle bounds from the ACME II result. Fig. from Matt Reece (unpublished). ACME III projection (~10<sup>-30</sup> e cm) dashed. Key EDM results since 2010. Two-loop sensitivity from Nakai & Reece (2017). One-loop sensitivity from Feng (2013). LHC scale gives stop mass sensitivity.

# SECTION II: HOW CAN WE MEASURE AN ELECTRON EDM?

#### HOW DO WE MEASURE AN EDM?

An eEDM will precess in both electric and magnetic fields

$$\begin{split} E &= -\vec{d_e} \cdot \vec{\mathcal{E}} - \vec{\mu} \cdot \vec{B} \\ \phi(\vec{\mathcal{E}}) &= E\tau/\hbar = -(\vec{d_e} \cdot \vec{\mathcal{E}} + \vec{\mu} \cdot \vec{B})\tau/\hbar \end{split} \stackrel{\vec{\mathbf{E}}_{\text{eff}} \ \vec{B}}{\uparrow} \quad \hat{\vec{x}} \stackrel{\hat{\vec{\mathcal{E}}}_{\text{eff}} \ \vec{B}}{\downarrow} \quad \hat{\vec{x}} \stackrel{\hat{\vec{\mathcal{E}}}}{\downarrow} \quad \hat{\vec{X}} \stackrel{\hat{\vec{\mathcal{E}}}}{\downarrow} \quad \hat{\vec{X}} \stackrel{\hat{\vec{\mathcal{E}}}_{\text{eff}} \ \hat{\vec{B}}}{\downarrow} \quad \hat{\vec{X}} \stackrel{\hat{\vec{\mathcal{E}}}_{\text{eff}} \ \hat{\vec{B}}}{\downarrow} \quad \hat{\vec{X}} \stackrel{\hat{\vec{\mathcal{E}}}}{\downarrow} \quad \hat{\vec{X}} \stackrel{\hat{\vec{B}}}{\downarrow} \quad \hat{\vec{X}} \stackrel{\hat{\vec{A}}}{\downarrow} \quad \hat{\vec{X}} \stackrel{\hat{\vec{A}}}{\vec{X}} \stackrel{\hat{\vec{A}}}}{\vec{X}} \stackrel{\hat{\vec{A}}}{\vec{X}} \stackrel{\hat{\vec{A}}}{\vec{X}} \stackrel{\hat{\vec{A}}}{\vec{X}} \stackrel{\hat{\vec{A}}}{\vec{X}} \stackrel{\hat{\vec{A}}}{\vec{X}} \stackrel{\hat{\vec{A}}}{\vec{X}} \stackrel{\hat{\vec{A}}}{\vec{X}} \stackrel{$$

Reverse E field for differential measurement

$$\phi_{EDM} = (\phi(\vec{\mathcal{E}}) - \phi(-\vec{\mathcal{E}}))/2$$
$$\phi_{EDM} = -(\vec{d_e} \cdot \vec{\mathcal{E}})\tau/\hbar$$



### AN EDM IN THO

- Molecules can provide strong internal electric fields
  - $\blacktriangleright$  ThO has  $E_{eff} \sim 80 \; GV/cm$
- We can flip this electric field by probing different states
  - Powerful method for eliminating systematic errors
  - ThO only requires a small applied field





 $|m_{\mathbf{J}}=0\rangle$ 



### **MOLECULAR EDMS**

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# **EXPERIMENTAL SENSITIVITY**

For a shot noise limited measurement we

$$\delta\omega \propto \frac{1}{\tau\sqrt{\dot{N}T}}$$
  $\tau$  Coherence time  
 $\dot{N} = \text{Count rate}$   
 $T = \text{Averaging time}$ 

For our experiment we have  $\mathcal{H}_{EDM} = -\vec{d_e} \cdot \vec{\mathcal{E}_{eff}}$ 

$$\delta d_e \propto \frac{1}{\mathcal{E}_{eff} \tau \sqrt{\dot{N}T}}$$
  $\mathcal{E}_{eff} = \text{Electric field}$ 

#### **EXPERIMENT STRUCTURE**







# **STATE PREPARATION**

- Stimulated Raman Adiabatic
  Passage (STIRAP)
  - Coherent population transfer from X to H
  - ~75% transfer efficiency
- State Refinement
  - Optically pump into dark state with desired polarization
  - Suppress residual STIRAP phases



# **STATE READOUT**

- Project phase onto orthogonal polarizations
- Rapidly switch polarization at 200 kHz
  - AOMs allow rapid switching
- Detect fluorescence with 8 PMTs





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#### **ACME II RESULT**

ACME Collaboration et al., Nature (2018):

$$d_e = (4.3 \pm 3.1_{stat} \pm 2.6_{syst}) \times 10^{-30} e \text{ cm}$$

Implies a 90% C. L. upper limit of:

 $|d_e| < 1.1 \times 10^{-29} e \text{ cm}$ 

Probes new physics up to ~30
 TeV at 1-loop level

Table 1   Systematic shifts for $\omega^{NC}$ and	their statisti	cal uncertainties
Parameter	Shift	Uncertainty
$\partial \mathcal{B}_z / \partial z$ and $\partial \mathcal{B}_z / \partial y$	7	59
$\omega_{\rm ST}^{NE}$ (via $\theta_{\rm ST}^{\rm H-C}$ )	0	1
$P_{ref}^{\mathcal{NE}}$	-	109
$\mathcal{E}^{nr}$	-56	140
$\mathcal{C} ^{\mathcal{NE}}$ and $ \mathcal{C} ^{\mathcal{NEB}}$	77	125
$\omega^{\mathcal{E}}$ (via $\mathcal{B}_{z}^{\mathcal{E}}$ )	1	1
Other magnetic-field gradients (4)	-	134
Non-reversing magnetic field, $\mathcal{B}_z^{nr}$	-	106
Transverse magnetic fields, $\mathcal{B}_x^{nr}$ , $\mathcal{B}_y^{nr}$	-	92
Refinement- and readout-laser detunings	-	76
$ ilde{\mathcal{N}}$ -correlated laser detuning, $ extsf{D}^{\mathcal{N}}$	-	48
Total systematic	29	310
Statistical uncertainty		373
Total uncertainty		486

Values are shown in  $\mu$ rad s<sup>-1</sup>. All uncertainties are added in quadrature. For  $\mathcal{E}_{eff} = 78 \text{ GV cm}^{-1}$ ,  $d_e = 10^{-30}e \text{ cm}$  corresponds to  $|\omega^{\mathcal{N}\mathcal{E}}| = \mathcal{E}_{eff}d_e/\hbar = 119 \,\mu$ rad s<sup>-1</sup>.

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# SECTION III: A NEW ACME MEASUREMENT

### ACME III OVERVIEW



# **INCREASED H-STATE LIFETIME**

- In January measured H-state lifetime to be 5.3 ms
- Last lifetime measurement showed a lower bound of just 1.8 ms
- ACME II used only 1 ms (20 cm) precession time
- Currently working on measurement to reduce uncertainty of H-state lifetime



#### **ELECTROSTATIC LENS FOR THO MOLECULES**

- Without a lens, fewer than .04% of molecules reach the detector
- Electrostatic lens focuses molecules into the EDM region, giving
  ~20x gain in signal (including the efficiency of double-STIRAP)
- Efficient STIRAP into the Q state allows for strong focusing





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## CONTROLLING EXCESS NOISE FOR ACME III

- ACME II had 1.7 times more noise than expected from the shot noise limit
- Noise came from 2 effects:
  - Large scale timing jitter
  - Timing offset between X and Y polarization bins
  - We can now control both parameters to reduce this noise





# **DETECTION UPGRADES**

- SiPMs provide a significant gain over the PMTs used in ACME II
- Increased light pipe diameter improves our geometric collection efficiency

	Requirement	Measured	Comment
Photon detection efficiency	~ 50%	$PMT \times 2.5$	Absolute value is unknown.
Dark count rate	< 10 Mcps	< 10 Mcps	Cooled down to -10°C
Cross talk & After pulse	< 25%	~ 20%	Array type package.
3dB Bandwidth	5 MHz	8.5 MHz	w/ Pole-Zero Cancellation
Electrical noise	< 10 nV/√Hz	< 10 nV/√Hz	

## ACME III PROPOSED GAINS

$$\delta d_e = \frac{1}{2T \mathcal{E}_{eff} \sqrt{N}}$$

Improvement	Signal Gain	EDM Sensitivity Gain
ncreased Precession Time	0.20	2.3
Electrostatic Lens	20.5	4.5
SiPM Detector Upgrade	2.3	1.5
Timing Jitter Noise Reduction	1	1.7
Total	9.4	26.4

#### OUTLOOK

- EDM searches are a powerful tool for studying BSM physics
- The ACME II measurement placed the strongest limit on the eEDM to date
- Promising path towards a new measurement with an order of magnitude improvement in sensitivity

#### THE ACME COLLABORATION

#### Yale

David DeMille (PI) Xing Wu (postdoc) James Chow (grad student) Zhen Han (grad student) Peiran Hu (grad student)

#### Harvard

John Doyle (PI) Xing Wu (postdoc)

#### Northwestern

Gerald Gabrielse (PI) Daniel Lascar (Research Asst. Prof.) Daniel Ang (Harvard grad student) Siyuan Liu (grad student) Bingjie Hao (grad student)

#### **Okayama University**

Takahiko Masuda (Associate Professor) Koji Yoshimura (Professor) Satoshi Uetake (Associate professor) Noboru Sasao (Professor)





James Chow Zhen Han

Dan Lascar



Xing Wu



John Doyle

David DeMille



Gerald Gabrielse



Siyuan Liu



Cole M.



Bingjie Hao



Daniel Ang